LETTER

Two-parameter areal scaling in the Hénon map

To cite this article: Gonzalo Marcelo Ramírez-Ávila et al 2019 EPL 126 20001

View the article online for updates and enhancements.
Two-parameter areal scaling in the Hénon map

GONZALO MARCELO RAMÍREZ-ÁVILA\textsuperscript{1,2,3}, IMRE M. JÁNOSI\textsuperscript{1,3,4} and JASON A. C. GALLAS\textsuperscript{1,2,3,5}

\textsuperscript{1} Max-Planck Institute for the Physics of Complex Systems - Nöthnitzer Str. 38, 01187 Dresden, Germany
\textsuperscript{2} Instituto de Investigaciones Físicas, Campus Cota-Cota, Universidad Mayor de San Andrés - La Paz, Bolivia
\textsuperscript{3} Instituto de Altos Estudios da Paraíba - Rua Silvino Lopes 419-2502, 58039-190 João Pessoa, Brazil
\textsuperscript{4} Department of Physics of Complex Systems, Eötvös Loránd University - H-1117 Budapest, Hungary
\textsuperscript{5} Complexity Sciences Center - 9225 Collins Avenue Suite 1208, Surfside, FL 33154, USA

received 21 February 2019; accepted in final form 15 April 2019
published online 29 May 2019

PACS 02.70.-c – Computational techniques; simulations
PACS 02.70.Rr – General statistical methods

Abstract – We study a bifurcation cascade whose proper unfolding requires tuning more than one parameter simultaneously. Specifically, we investigate metric properties of extended self-similar triangular areas observed recently in the control parameter space of flows (lasers and electronic circuits), and maps. Such areas are delimited by shrimplike stability islands, seem to arise in unbounded quantities, and to accumulate in narrow intervals of control parameters. Numerically, we find their asymptotic rate of accumulation to be unity. The asymptotic properties of triangle vertices and their centroids are also investigated.

Copyright © EPLA, 2019

Introduction. – Recently, a profusion of zig-zag networks interconnecting certain classes of periodic oscillations were discovered in the control parameter space of a fiber-ring laser, in an electronic circuit containing a tunnel diode [1,2], and in the Hénon map, a proxy for a widely used class of CO\textsubscript{2} lasers [3,4]. Zig-zag networks consist of regular chains interconnecting sequences of intricate and self-similar stability phases known as shrimps [5–10], formed by pairs of cascades of either period or peak doubling bifurcations followed by chaotic oscillations. Such networks are not difficult to find in both continuous-time and discrete-time dynamical systems.

One of the distinctive characteristics of zig-zag networks is that they sometimes display infinite accumulation of shrimp triplets which form triangles, as illustrated below. Thus, they offer a natural scenario to investigate metric properties of the accumulation of bifurcation cascade whose proper unfolding requires tuning simultaneously more than one parameter. In particular, zig-zag networks allow the investigation of scaling properties of extended areas discovered recently in the control parameter space of prototypical systems, namely in the self-pulsations of a CO\textsubscript{2} laser with feedback [10,11], in a damped-driven Duffing oscillator [12], and in the characterization of the transport properties of ratchets [13–15]. Accordingly, the present work grew out of a desire to study scaling properties of stability islands whose generic shape and position in control parameter space depend on tuning more than one control parameter simultaneously. Multiparameter scalings do not seem to have been explored yet.

As is known, the investigation of metric properties of bifurcation cascades was the subject of several studies probing universality classes in dynamical systems. Such studies were motivated originally by remarkable findings reported independently by Feigenbaum [16] and by Coullet and Tresser [17,18]. For more recent results see, e.g., refs. [19,20]. Despite the initial claims of universality of the scaling constants, it was concomitantly reported by several groups that the scaling constants, in fact, vary considerably in systems more complex than the quadratic map, and in higher dimensions [21–31].

Concerning metric properties, period-doubling bifurcations in low-dimensional systems have been studied extensively. However, such investigations were restricted exclusively to properties observed when varying a single control parameter. As is known, the most pronounced effects of bifurcation cascades occur along certain specific directions, tortuous corridors in parameter space, which invariably require tuning more than one parameter in order to be able to move along them [5,10]. Here, we focus on metric properties observed when complex extended structures in parameter space are deformed by the simultaneous variation of two control parameters. Clearly, the need for tuning more than one parameter simultaneously arises...
investment of computer time. A significant advantage of numerical systems. Accordingly, such investigations must be performed numerically. For practical applications, the identification of complex structures and their accumulation mechanisms in maps can be made with a moderate investment of computer time. A significant advantage of studying metric properties of maps is the possibility to bypass all the usual uncertainties associated with numerical algorithms used for the integration of sets of differential equations.

**Shrimp doublets and triplets.** – Figure 1 shows a broad view of the control parameter space of the Hénon map, the region where one finds most of the shrimp-like islands of stability [5–8]. Numbers indicate the main period \( k \) of some of the \( k \times 2^n \) islands.

Rather than using eigenvalues [32], in fig. 1 we follow Sannami [33] in plotting the trace \( \tau_k \) of the Jacobian matrix for \( k \)-periodic points. The reason for using the trace is that eigenvalues are not always real numbers and have manifolds that may display odd behaviors [34]. Therefore, eigenvalues do not seem reliable to inspect the inner structure of shrimps. Instead of using a single solid color to paint the whole \( k \)-periodic phase, we partitioned phases into two colored sectors as follows. For a given period \( k \), we represented the region characterized by \( \tau_k > 0 \) using a color associated with the period, using black to represent the region where \( \tau_k < 0 \). This dichotomic division of the stability windows, the same used in all figures below, increases the information content of stability diagrams, allowing one to easily recognize shrimps sharing similar periodicities and, simultaneously, revealing their inner structure, analogously to plots of “multipliers” for one-dimensional maps [35].

In fig. 1, the white region represents parameters leading to aperiodic (i.e., chaotic) orbits. Starting from the left side, fig. 1 shows two pairs of stripes containing periods 2 and 4, as indicated. They belong to the familiar \( 1 \times 2^n \) bifurcation cascade. After the rightmost period-4 region, it is possible to recognize a similar pair of parabolic stripes corresponding to period 8, also characterized by negative and positive values of \( \tau \). In the upper part of the period-8 cascade, there is a black box containing a large portion of an additional complicated period-8 structure, which extends well into the vast parameter region characterized by divergence, as indicated. This additional period-8 island contains a cusp located somewhat near structures of periods 10 and 6. Incidentally, around these islands one finds a startling phenomenon: *stable periodic orbits characterized by complex values of \( (x, y) \) but for real parameters \((a, b)\) [36].

Figure 1 also contains two boxes with shrimp doublets and triplets. As mentioned, the large and easily visible box contains part of the period-8 structure. A second and much smaller box, indicated by an arrowhead, is located between shrimps of periods 7 and 9. It is shown magnified in fig. 2(a). At the center of this figure there is a wide period-18 stability island mentioned by Lorenz [7]. As is clear from the figure, the trace \( \tau \) reveals a relatively complicated inner topography of the central portion of this island. On a finer scale, around the period-18 island there is a profusion of shrimp doublets, some of which are shown in fig. 2(b). Sometimes, such doublets are in fact triplets.
Two-parameter areal scaling in the Hénon map

Fig. 2: Sequences of shrimp doublets and triplets. Numbers indicate the period of the main stability region. The white background represents parameters leading to chaotic oscillations. The pink background is the basin of the attractor at $-\infty$.

(a) The complex period-18 structure studied by Lorenz [7], surrounded by shrimp doublets and triplets. Boxes are magnified in the next three panels. (b) A sequence of shrimp doublets. (c) An apparently isolated pair of period-25 shrimps which, however, forms (d) a shrimp triplet; (e) a period-23 triplet between a pair of period-29 triplets; (f) a region with a profusion of triplets and more intricate stability islands. The two boxes are magnified in fig. 6. Individual panels display the analysis of $1200 \times 1200 = 1.44 \times 10^6$ parameter points.

Fig. 3: Successive magnifications illustrating a profusion of zig-zag triplets. (a) Magnification of the uppermost box in fig. 1. (b) Enlargement of the red box in (a). (c) Enlargement of the black box in (a). (d) Apparently an unbounded arithmetic progression of zig-zag triplets is located in the box, shown magnified in fig. 5. Similar sequences exist in other regions of the control parameter space. Grid resolution: $1200 \times 1200$ parameter points.
which may also arise as combinations of unsuspected and apparently uncorrelated structures, as shown in fig. 2(c) and fig. 2(d).

As hinted by the periodicities of individual doublets in fig. 2(b), they do not seem to be connected in any noticeable way. Uncorrelated doublets exist also in several other locations in the control parameter plane. Analogously, there is a large number of triplets, like in fig. 2(d), which do not seem to be connected to other stability islands. Attempts to detect shrimp connections met difficulties because their legs get thinner and thinner as one moves away from their central stability region. Similarly to fig. 2(d), fig. 2(e) illustrates a period-23 triplet formation in the same parameter region where there are two period-29 triplets. Such mixed formations are also found in other windows in the $b > 0$ half of the control plane. Figure 2(f) shows a sort of “border line” triplet straddling the chaotic and the divergent backgrounds, namely a triplet having two shrimps located over a background of chaos linked to a shrimp located over the background of divergence. Near this triplet, one finds a plethora of additional triplets as well as more complicated arrangements, illustrated by the pair of boxes in fig. 2(f), shown magnified below in fig. 6. In contrast to the isolated doublets and triplets in fig. 2, it is also possible to find unbounded cascades of self-similar triplets forming arithmetic progressions, namely whose periodicity increases by a constant value from triplet to triplet, as illustrated in fig. 5 and discussed in the next section.

**Triplets in arithmetic progression.** – Figure 3 shows a sequence of successively magnified windows indicating the location of an interesting arithmetic progression of shrimp triplets that we wish to consider in more detail. The pair of boxes fig. 3(a) contains several triplets analogous to the ones observed in systems governed by differential equations, namely in fiber-ring lasers, and in an electronic circuit with a tunnel diode [1,2]. Figure 3(b) shows uncorrelated triplets similar to the ones in fig. 2(d), while the red boxes in figs. 3(c) and (d) mark the location of triplets in an apparently never-ending arithmetic progression. Similar unbounded progressions exist in other parameter windows, particularly for orbits of higher periods. Such apparently unbounded progressions of stability islands display accumulation boundaries, horizons, embedded in the broad parameter background associated with chaotic oscillations.

The study of metric properties of the two-dimensional Hénon map and higher-dimensional maps is more complicated than the corresponding study for one-dimensional systems. For one-dimensional maps $x_{i+1} = f(x_i)$, the study of metric properties is greatly facilitated by the presence of critical points, namely points where $df(x)/dx|_{x=x_0} = 0$. Orbits containing such points are the so-called superstable orbits. For such orbits, the multiplier $m_k \equiv df_k/dx$ associated with a $k$-periodic orbit is zero [37–39] ($f_k$ denotes the $k$-th composition of $f$ with itself) [35]. Critical points are the basic objects used by Fatou and Julia to study the properties of iterated rational functions. For a very complete survey of the classical literature see Cremer [39]. For more recent literature consult ref. [40]. Unfortunately, for high-dimensional maps there are no proper definitions for critical points, multipliers, and superstable orbits.

Figure 4 shows enlarged views of the three shrimps forming the vertices $A_1B_1C_1$ and $A_2B_2C_2$ of the first two triplets in arithmetic progression. From fig. 4 one clearly sees that the trace of the Jacobian matrix is not equivalent to the multiplier. For, although the trace is capable of exposing two parabolic arcs which resemble the parabolas generated by multipliers for one-dimensional maps, for the Hénon map the parabolic arcs are “broken”, i.e., they do not always intersect, as in panels $A_1, B_1, A_2, B_2, C_2$. Furthermore, when they do intersect, the intersection occurs not at just a single point but, instead, in an extended region, as seen in panel $C_1$. These two problems are generic difficulties present in all higher-dimensional systems. To
Two-parameter areal scaling in the Hénon map

Table 1: Period $k_i$, coordinates and centroids of the triangles in arithmetic progression, shown in fig. 5. The values in the bottom line are extrapolated values. See text.

<table>
<thead>
<tr>
<th>$i$</th>
<th>$k_i$</th>
<th>$a_{A_i}$</th>
<th>$b_{A_i}$</th>
<th>$a_{B_i}$</th>
<th>$b_{B_i}$</th>
<th>$a_{C_i}$</th>
<th>$b_{C_i}$</th>
<th>Area $\times 10^6$</th>
<th>$a_{centroid}$</th>
<th>$b_{centroid}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>44</td>
<td>1.12878432</td>
<td>0.42556190</td>
<td>1.1292456</td>
<td>0.42542796</td>
<td>1.12917220</td>
<td>0.42507113</td>
<td>5.95696512</td>
<td>0.12906703</td>
<td>0.42556666</td>
</tr>
<tr>
<td>2</td>
<td>62</td>
<td>1.12942445</td>
<td>0.42515476</td>
<td>1.12958722</td>
<td>0.42513898</td>
<td>1.12966421</td>
<td>0.42525273</td>
<td>0.98757841</td>
<td>0.12955386</td>
<td>0.42518213</td>
</tr>
<tr>
<td>3</td>
<td>80</td>
<td>1.12957064</td>
<td>0.42504772</td>
<td>1.12967100</td>
<td>0.42505563</td>
<td>1.12976177</td>
<td>0.42513312</td>
<td>0.36404728</td>
<td>0.12966067</td>
<td>0.42507882</td>
</tr>
<tr>
<td>4</td>
<td>98</td>
<td>1.12962318</td>
<td>0.42508662</td>
<td>1.12970500</td>
<td>0.42502402</td>
<td>1.12979458</td>
<td>0.42508885</td>
<td>0.19603979</td>
<td>0.12970762</td>
<td>0.42504066</td>
</tr>
<tr>
<td>5</td>
<td>116</td>
<td>1.12964741</td>
<td>0.42499058</td>
<td>1.12971932</td>
<td>0.42500938</td>
<td>1.12980916</td>
<td>0.42506822</td>
<td>0.12710962</td>
<td>0.12972530</td>
<td>0.42502737</td>
</tr>
<tr>
<td>6</td>
<td>134</td>
<td>1.12966058</td>
<td>0.42498084</td>
<td>1.12972700</td>
<td>0.42500162</td>
<td>1.12981689</td>
<td>0.42505705</td>
<td>0.90608732</td>
<td>0.12973482</td>
<td>0.42503171</td>
</tr>
<tr>
<td>7</td>
<td>152</td>
<td>1.12966845</td>
<td>0.42497097</td>
<td>1.12973143</td>
<td>0.42499706</td>
<td>1.12982147</td>
<td>0.42505030</td>
<td>0.06820358</td>
<td>0.12974045</td>
<td>0.42500744</td>
</tr>
<tr>
<td>8</td>
<td>170</td>
<td>1.12967360</td>
<td>0.42497116</td>
<td>1.12973422</td>
<td>0.42499418</td>
<td>1.12982446</td>
<td>0.42504592</td>
<td>0.05295770</td>
<td>0.12974049</td>
<td>0.42500375</td>
</tr>
<tr>
<td>9</td>
<td>188</td>
<td>1.12967713</td>
<td>0.42496655</td>
<td>1.12973609</td>
<td>0.42499224</td>
<td>1.12983644</td>
<td>0.42504289</td>
<td>0.04295662</td>
<td>0.12974655</td>
<td>0.42500123</td>
</tr>
<tr>
<td>10</td>
<td>206</td>
<td>1.12967971</td>
<td>0.42496070</td>
<td>1.12973741</td>
<td>0.42499888</td>
<td>1.12983879</td>
<td>0.42504073</td>
<td>0.03442693</td>
<td>0.12974634</td>
<td>0.42500441</td>
</tr>
<tr>
<td>11</td>
<td>224</td>
<td>1.12968156</td>
<td>0.42496529</td>
<td>1.12973837</td>
<td>0.42498889</td>
<td>1.12983896</td>
<td>0.42503912</td>
<td>0.02841212</td>
<td>0.12974963</td>
<td>0.42499810</td>
</tr>
<tr>
<td>55</td>
<td>1016</td>
<td>1.12968365</td>
<td>0.42496342</td>
<td>1.12973944</td>
<td>0.42498857</td>
<td>1.12982988</td>
<td>0.42503696</td>
<td>0.02129941</td>
<td>0.12975092</td>
<td>0.42499632</td>
</tr>
</tbody>
</table>

Fig. 5: The first 11 triangles of an apparently infinite arithmetic progression accumulating towards a period-18 boundary. The difference of the periods between two consecutive triangles is also 18. Grid resolution: 3000 × 3000 parameter points.

bypass trace peculiarities and to be able to define unambiguously all shrimp heads [35], here we interpolated broken parabolic arcs and used their points of intersection to define triangle vertices.

Areal scaling. – Figure 5 shows the location and the strong compression undergone by the first 11 triangle triplets which accumulate in arithmetic progression towards the period-18 boundary as they successively get more and more squeezed. Red dots mark the centroid of the triangles, namely the intersection of the three triangle medians. The coordinates of the triangle vertices are recorded, their area, and their centroid coordinates are collected in table 1. These numerical values were obtained by measuring them from individual blowups (not given here) for every triangle. Noteworthy is the fact that the period difference between two consecutive triangles is 18, the same period boundary horizon towards which they accumulate. As mentioned above, this situation is analogous to the one previously observed in a damped-driven Duffing oscillator [12] and in the self-pulsations of a CO$_2$ laser with feedback [10,11].

As seen from fig. 5, vertices tend to accumulate fast, in a narrow parameter interval. This tendency may also be seen in table 1. Accordingly, an interesting issue is to determine their accumulation points and rate of convergence. To find them, we proceed as follows: i) Firstly, we compute the successive differences between the coordinates $(a, b)$ for vertices and centroids of each triangular region; ii) from these differences, a fitting equation for each sequence is derived; iii) using these fitting equations we estimate the coordinates for extrapolated triangles; iv) the extrapolation process is extended until quantities remain constant to eight decimal digits. The convergence rate of the triangles towards the asymptotic horizon is found to be unity. The resulting extrapolated values are listed in the last line of both tables above. For $i = 55$, the listed values for the area and centroid were obtained from the extrapolation, not from the vertices coordinates in the table, although both sets of values essentially coincide. Remarkably, triangles seem to accumulate just before reaching the period-18 horizon leg in front of them. Perhaps extrapolations using more than 11 triangles could reveal the extrapolated values coming closer or even coinciding asymptotically with the convergence horizon. However, it becomes increasingly more difficult to reliably detect triangles when the period further increases. The precise location of the convergence horizon is therefore left as an open question for further investigation.

As a last result, in fig. 6 we collect a number of triangular stability islands which “break the symmetry”, namely that do not fit unambiguously in the above scenarios but, instead, display more exquisite shapes and organizations. For instance, the box in fig. 6(a), magnified in fig. 6(c), displays a pair of shrimps that under low resolution may appear as uncorrelated but that are in fact interconnected, forming a triplet. Figure 6(d) shows that a period-22 triplet is partially overlapping the leftmost partner of a larger period-16 triplet. In reality, the period-22 triplet is interconnected with a fourth shrimp located farther to the right, as indicated. Therefore, it is possible to circulate continuously from one shrimp to the others without ever having to cross the vast sea of chaos surrounding
Fig. 6: Magnifications of the pair of boxes seen in fig. 2(f) showing some exquisite and much more complicated triplets.

Conclusions and outlook. – We studied metric properties of certain triangular stability islands covering extended areas in control parameter space and which are abundantly present in flows and maps. Such triangular islands appear both as isolated forms or as forming apparently unbounded arithmetic progressions. In contrast with the familiar scalings in the literature, the unfolding of areal scaling requires tuning more than one control parameter simultaneously. A significant feature of the arithmetic progression is that it displays specific accumulation points, both for triangle vertices and their centroid coordinates. Although the emphasis here was on a specific period-18 accumulation, we find accumulations to be a rather common phenomenon, involving analogous arithmetic progressions and many other periods. The accumulation unfolds systematically and is fast. Accordingly, the convergence to an almost constant value of the area was observed. Furthermore, we find the arithmetic progression to converge to a well-defined asymptotic horizon whose period coincides with the constant rate of period increase of the arithmetic progression. It is not yet clear if the arithmetic progressions involve a finite or an infinite number of terms. A particularly promising system for investigating two-parameter scalings is the analytical path discussed in fig. 2 of ref. [41], for the so-called canonical quartic map. In conclusion, the metric properties of extended progressions of stability structures whose accumulation in control parameter space depends on more than one parameter were studied in detail and characterized numerically. We are not aware of any previous study of the scaling of properties depending on the variation of more than one parameter simultaneously. Our results are also relevant for flows, systems governed by differential equations. It would be interesting to compare the present findings with analogous ones for ratchets and the aforementioned flows representing semiconductor laser diodes, electronic circuits, and other promising systems.

***

This work was supported by the Max-Planck Institute for the Physics of Complex Systems, Dresden, in the framework of the Advanced Study Group: Forecasting with Lyapunov vectors. IMJ was supported by the Hungarian National Research, Development and Innovation Office under grant K-125171. JACG was also supported by CNPq, Brazil.

REFERENCES

Two-parameter areal scaling in the Hénon map

[1] Two-parameter areal scaling in the Hénon map


